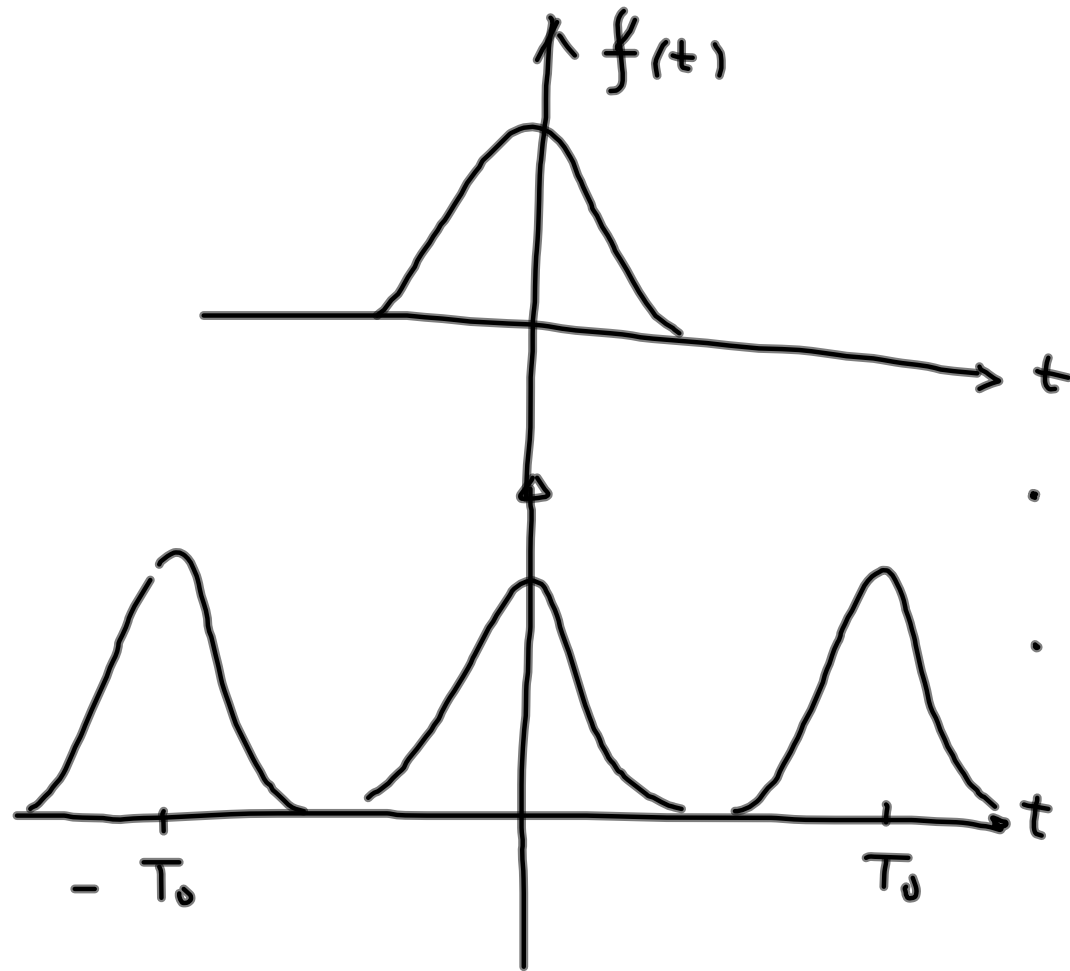


Chapter 3. Fourier Transform & Frequency-domain Analysis

§ 3.1. Fourier Transform

① Aperiodic Signal Representation by Fourier Integral



- Aperiodic signal $f(t)$
- Construct a periodic signal $f_{T_0}(t)$ by repeating $f(t)$ every T_0
- Let $T_0 \rightarrow \infty$
 $\lim_{T_0 \rightarrow \infty} f_{T_0}(t) = f(t)$

Fourier Series of $f_T(t)$ ω

$$f_T(t) = \sum_{n=-\infty}^{+\infty} D_n \underbrace{e^{j n \omega_0 t}}_{\omega_0 = \frac{2\pi}{T_0}}$$

$$D_n = \frac{1}{T_0} \int_{-T_0/2}^{T_0/2} f_T(t) e^{-j n \omega_0 t} dt$$

as $T_0 \rightarrow \infty$:

$$\omega_0 = \frac{2\pi}{T_0} \rightarrow 0 \quad \omega_0 \leftarrow \Delta\omega$$

$$\therefore \int_{-T_0/2}^{T_0/2} f_T(t) e^{-j n \omega_0 t} dt$$

$$\rightarrow \int_{-\infty}^{\infty} f(t) e^{-j \boxed{n \Delta\omega} t} dt = F(n\Delta\omega)$$

where $F(\omega) \stackrel{\text{def}}{=} \int_{-\infty}^{+\infty} f(t) e^{-j\omega t} dt$

$$\Rightarrow D_n = \frac{1}{T_0} F(n \cdot \Delta\omega) = \frac{\Delta\omega}{2\pi} \cdot F(n \cdot \Delta\omega)$$

$$\therefore f_{T_0}(t) = \sum_{h=-\infty}^{+\infty} \underbrace{\frac{F(n \cdot \Delta\omega) \cdot \Delta\omega}{2\pi}}_{D_n} e^{jn \cdot \Delta\omega \cdot t}$$

$$\Rightarrow \underline{\underline{f(t)}} = \lim_{T_0 \rightarrow \infty} f_{T_0}(t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \underbrace{F(\omega)}_{\substack{\text{red arrow} \\ \text{to } D_n}} \underbrace{e^{j\omega t}}_{\text{red box}} d\omega$$

Fourier transform:

$$F(\omega) = \mathcal{F}\{f(t)\} = \int_{-\infty}^{+\infty} f(t) e^{-j\omega t} dt$$

Inverse Fourier transform

$$f(t) = \mathcal{F}^{-1}\{F(\omega)\} = \frac{1}{2\pi} \int_{-\infty}^{+\infty} F(\omega) e^{j\omega t} d\omega$$

(2) Sufficient conditions for existence
of Fourier transform

• Weak condition:

$$\text{let } \hat{f}(t) \stackrel{\text{def}}{=} \frac{1}{2\pi} \int_{-\infty}^{+\infty} F(\omega) e^{j\omega t} d\omega$$

$$e(t) = f(t) - \hat{f}(t)$$

Theorem 1:

If $f(t)$ has finite energy

$$\text{i.e., } \int_{-\infty}^{+\infty} |f(t)|^2 dt < \infty$$

then $\bar{F}(\omega)$ exists $\&$ $\int_{-\infty}^{+\infty} |k(t)|^2 dt = 0$.

$\rightarrow f(t) \& \hat{f}(t)$ may differ at discrete values of t . But the energy difference is 0.

• Strong condition



Theorem 2: If a) $\int_{-\infty}^{+\infty} |f(t)| dt < \infty$

Then $f(t) = \hat{f}(t)$ for every t , except at a discontinuity, where $\hat{f}(t)$ equal to the avg of the values on two sides of discontinuity.

b) $f(t)$ has a finite # of max & min within any finite interval.

c) $f(t)$ has a finite # of discontinuities within any interval; each of these discontinuities are finite

§ 3.2. Some useful Transforms

$$\textcircled{1} \quad \delta(t) \xleftrightarrow{\mathcal{F}} 1$$

$$\mathcal{F} \{ \delta(t) \} = \int_{-\infty}^{+\infty} \delta(t) e^{-j\omega t} dt = 1.$$

$$\textcircled{2} \quad 1 \xleftrightarrow{\mathcal{F}} 2\pi \delta(\omega)$$

$$\mathcal{F}^{-1} \{ 2\pi \cdot \delta(\omega) \} = 2\pi \cdot \frac{1}{2\pi} \int_{-\infty}^{+\infty} \delta(\omega) e^{j\omega t} d\omega = 1.$$

$$\textcircled{3} \quad e^{-at} u(t), \quad a > 0.$$

$$\xleftrightarrow{\mathcal{F}} \frac{1}{a+j\omega}$$

$$\mathcal{F} \{ e^{-at} u(t) \} = \int_{-\infty}^{+\infty} e^{-at} \underbrace{u(t)} e^{-j\omega t} dt$$

$$= \int_0^{\infty} e^{-(a+j\omega)t} dt$$

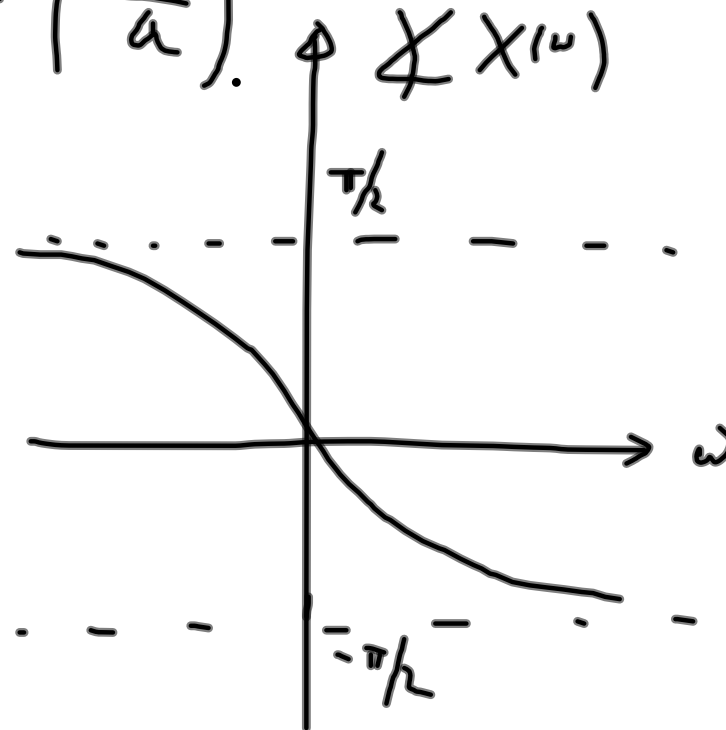
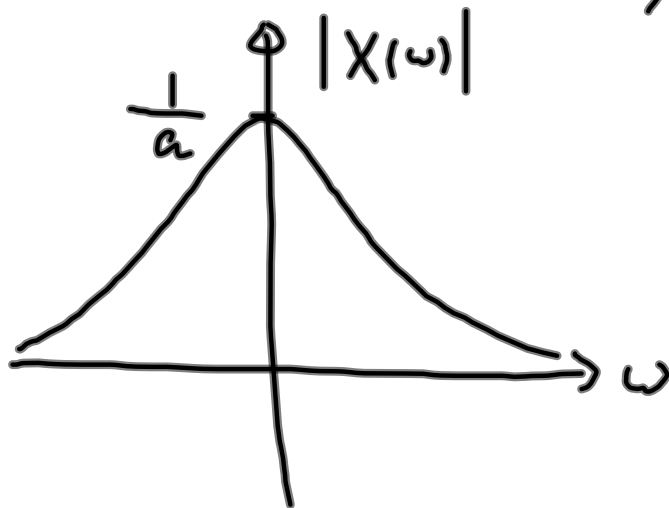
$$= -\frac{1}{a+j\omega} e^{-(a+j\omega)t} \Big|_0^{\infty}$$

$$= \frac{1}{a+j\omega}$$

$$X(\omega) = \frac{1}{a + j\omega} = \frac{a - j\omega}{a^2 + \omega^2}$$

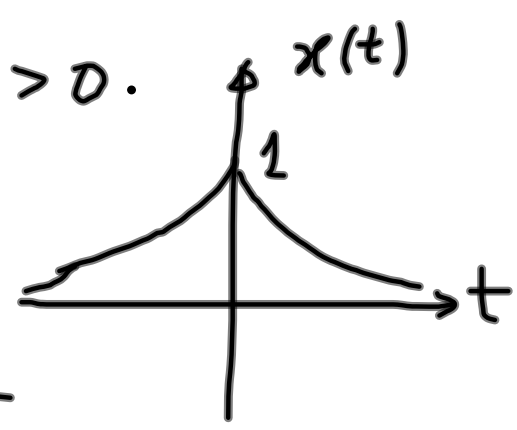
$$|X(\omega)| = \frac{1}{\sqrt{a^2 + \omega^2}}$$

$$\angle X(\omega) = -\tan^{-1}\left(\frac{\omega}{a}\right)$$



(4) $x(t) = e^{-a|t|} \quad a > 0.$

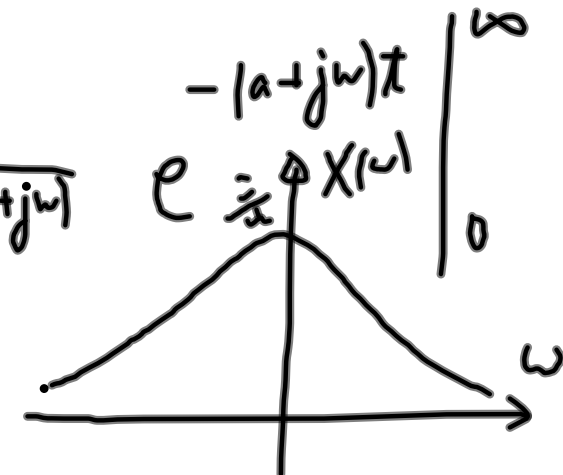
$\xleftrightarrow{f} \frac{2a}{\omega^2 + a^2}$



$$X(\omega) = \int_{-\infty}^{+\infty} e^{-a|t|} e^{-j\omega t} dt$$

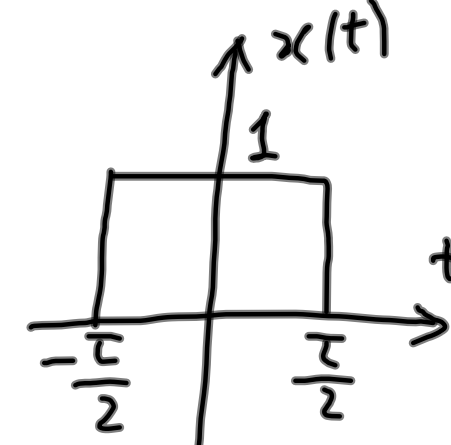
$$= \int_{-\infty}^0 e^{(a-j\omega)t} dt + \int_0^{\infty} e^{-(a+j\omega)t} dt$$

$$= \frac{1}{a-j\omega} e^{(a-j\omega)t} \Big|_{-\infty}^0 + \frac{1}{-(a+j\omega)} e^{-(a+j\omega)t} \Big|_0^{\infty}$$

$$= \frac{1}{a-j\omega} + \frac{1}{a+j\omega} = \frac{2a}{\omega^2 + a^2}$$


(5) $x(t) = \begin{cases} 1 & |t| < \frac{\tau}{2} \\ 0 & |t| > \frac{\tau}{2} \end{cases}$

$\xleftrightarrow{g} \tau \cdot \text{sinc}\left(\frac{\omega\tau}{2}\right)$



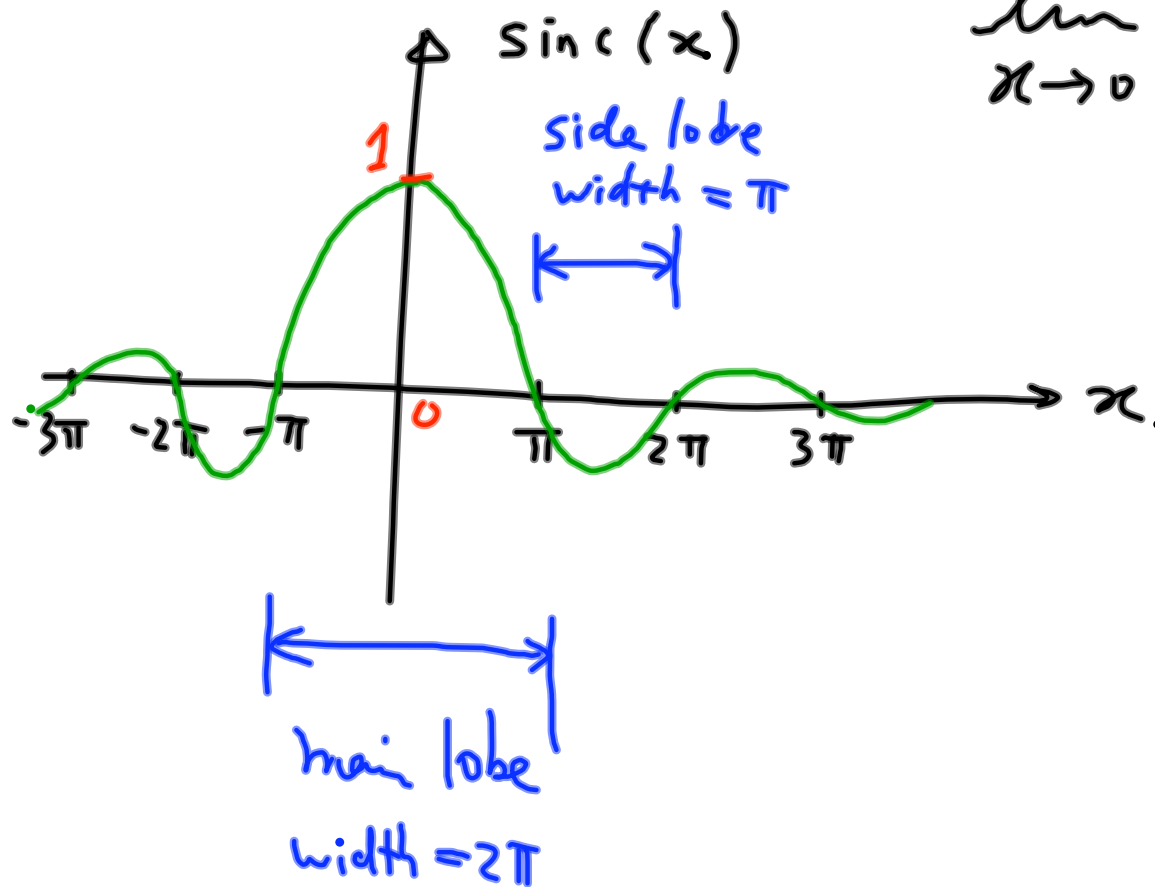
$$X(\omega) = \int_{-\frac{\tau}{2}}^{\frac{\tau}{2}} e^{-j\omega t} dt = \frac{1}{-j\omega} e^{-j\omega t} \Big|_{-\frac{\tau}{2}}^{\frac{\tau}{2}}$$

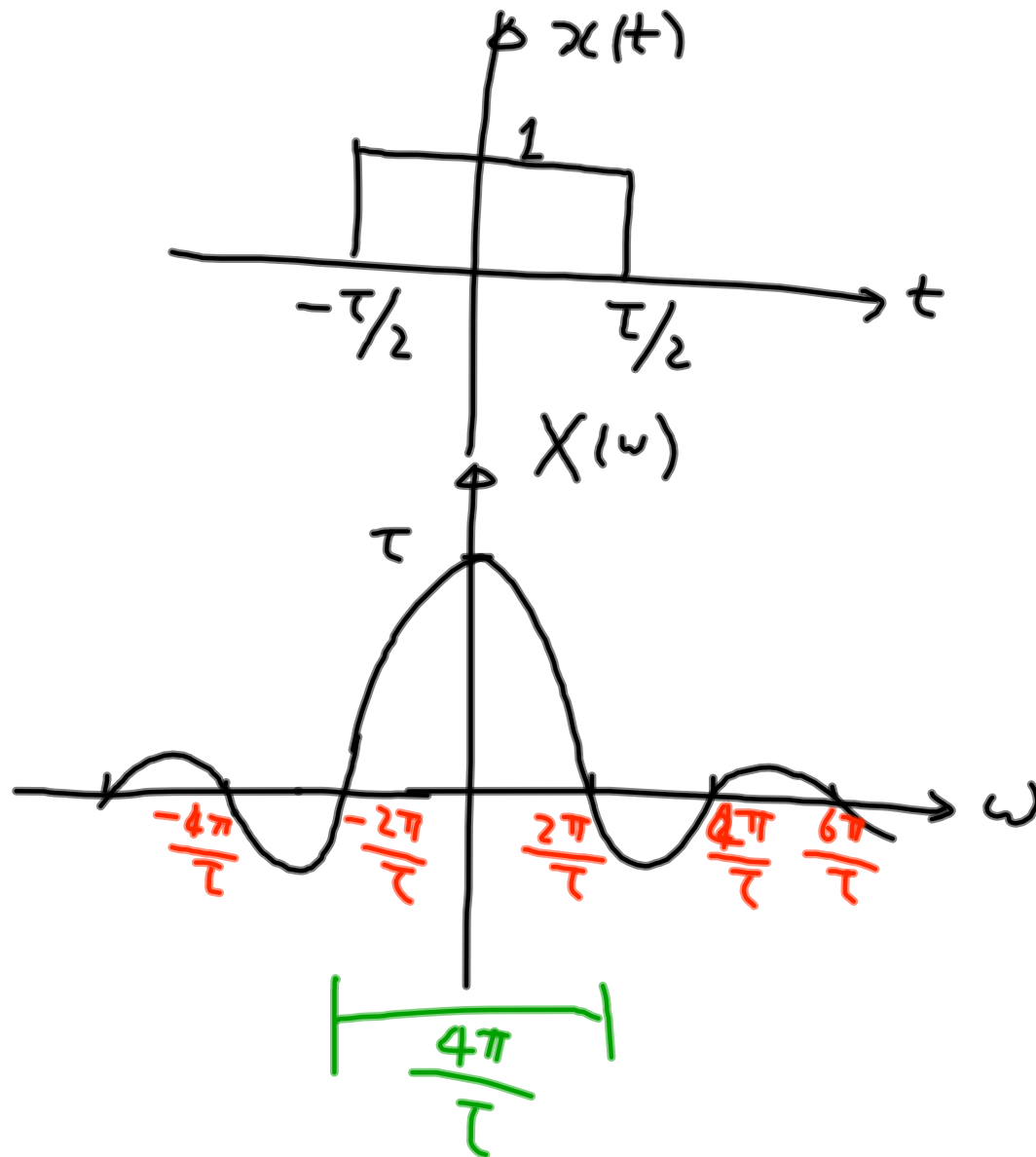
$$= \frac{e^{j\frac{\omega\tau}{2}} - e^{-j\frac{\omega\tau}{2}}}{j\omega} = \frac{2j \sin\left(\frac{\omega\tau}{2}\right)}{j\omega} = \tau \cdot \frac{\sin\left(\frac{\omega\tau}{2}\right)}{\left(\frac{\omega\tau}{2}\right)} = \tau \cdot \text{sinc}\left(\frac{\omega\tau}{2}\right)$$

$$\frac{\sin\left(\frac{\omega\tau}{2}\right)}{\left(\frac{\omega\tau}{2}\right)}$$

$$\text{sinc}(x) = \frac{\sin x}{x}$$

$$\lim_{x \rightarrow 0} \frac{\sin x}{x} = \frac{\cos x}{1} \Big|_{x=0} = 1.$$

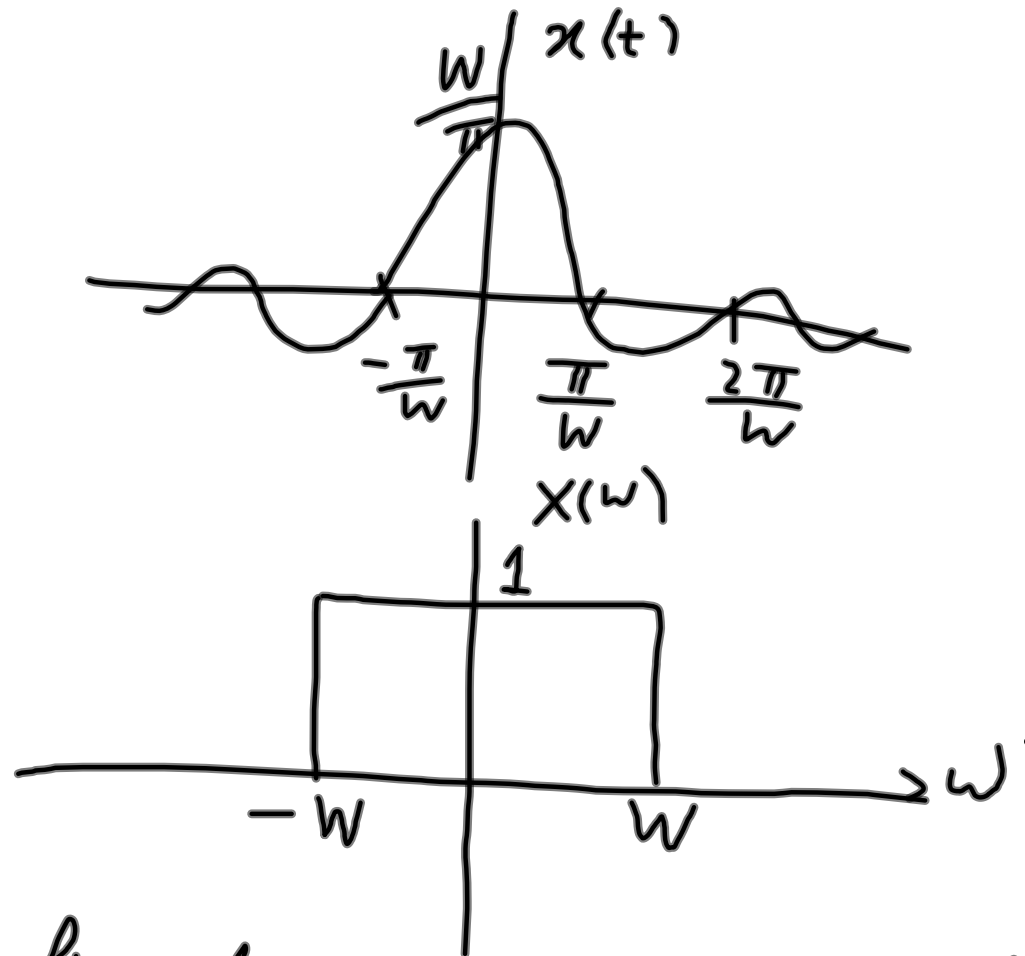




$$\textcircled{6} \quad x(t) = \frac{W}{\pi} \operatorname{sinc}(Wt)$$

$$\begin{array}{c} \mathcal{F} \\ \longleftrightarrow \end{array}
 X(\omega) = \begin{cases} 1 & |\omega| < W \\ 0 & |\omega| > W \end{cases}$$

$$\begin{aligned}
 x(t) &= \frac{1}{2\pi} \int_{-W}^W e^{j\omega t} d\omega \\
 &= \frac{\sin Wt}{\pi t} = \frac{W}{\pi} \operatorname{sinc}(Wt).
 \end{aligned}$$



freq-domain wider \leftrightarrow time-domain narrower